

## Hydrogen Bond Studies

### 67.\* The Crystal Structure of Rubidium Trihydrogen Selenite, $\text{RbH}_3(\text{SeO}_3)_2$

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The crystal structure of  $\text{RbH}_3(\text{SeO}_3)_2$  has been determined from three-dimensional single crystal X-ray diffractometer data obtained at room temperature. Four formula units crystallize in an orthorhombic unit cell of dimensions:  $a = 5.9192(2)$ ,  $b = 17.9506(5)$ , and  $c = 6.2519(3)$  Å. The space group is  $P2_1 2_1 2_1$ . The structure consists of two types of chains at a right angle. One chain is built up of  $\text{H}_2\text{SeO}_3$  molecules linked by 2.594(8)-Å hydrogen bonds and the other of  $\text{HSeO}_3^-$  ions linked by 2.571(12)-Å hydrogen bonds. These two types of chains are cross-linked by a third hydrogen bond of length 2.521(7) Å. The rubidium ion is surrounded by eight oxygen atoms forming a distorted cube. The  $\text{Rb}^+-\text{O}$  distances are in the range 2.94–3.19 Å.

#### Introduction

Crystals of the alkali trihydrogen selenites,  $\text{MH}_3(\text{SeO}_3)_2$ , have been found to show a variety of dielectric properties and crystal structures [for a review see Ref. (1)]. To be able to explain these different properties it is necessary to have detailed structural information.

The present investigation forms part of a series of X-ray and neutron diffraction studies of hydrogen-bonded ferroelectric and related compounds in progress at this Institute. The previous paper was a neutron diffraction study of  $\text{LiH}_3(\text{SeO}_3)_2$  (2), and the next will be an X-ray diffraction study of  $\text{NH}_4\text{H}_3(\text{SeO}_3)_2$  (3).

#### Experimental

Crystals of  $\text{RbH}_3(\text{SeO}_3)_2$  were prepared from an aqueous solution of  $\text{Rb}_2\text{CO}_3$  and  $\text{SeO}_2$  in the stoichiometric ratio. The solution was heated to 80°C. On cooling to room temperature small almost spherical single crystals were obtained. Their melting point was found to be 57°C, which agrees very well with the reported value of 57.1°C (1).

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The intensity data were collected at room temperature on a Stoe-Philips computer controlled four-circle diffractometer with a graphite monochromator using  $\text{MoK}\alpha$  radiation and  $\omega-2\theta$  scan. Different combinations of scan time and filter selection were used with the intention of measuring all reflections to about the same precision.

Three standard reflexions with integrated intensities of 3500, 4000, and 47 000 counts, respectively, were measured at regular intervals to provide a check on crystal and electronic stability. The two weaker reflexions showed no significant changes, whereas the intensity of the strongest one first underwent a continuous increase of 10% and then stabilized. All reflexions obtained up to this point were later remeasured. It was then found that the increase had affected only the high-intensity reflexions. This increase was probably due to an alteration in the nature of the extinction effects present. It could be that the mosaic structure had been changed because of radiation damage. As the data used in the refinements were those collected after the test reflexion had stabilized, no scaling was necessary.

A total of 1172 independent reflexions having  $\sin \theta/\lambda$  less than  $0.74 \text{ \AA}^{-1}$  were recorded. Out of these, 934 had intensities greater than  $3\sigma_{\text{count}}$ . The intensities were corrected for Lorentz, polarization, and absorption effects (spherical

crystal). The calculated linear absorption coefficient was  $189.7 \text{ cm}^{-1}$ . The crystal was approximately spherical with a radius of  $0.0376 \text{ mm}$ . The transmission factors were in the range  $0.356\text{--}0.373$ .

### Space Group and Unit Cell

The space group was determined from Weissenberg photographs recorded with  $\text{CuK}\alpha$  radiation. The diffraction symmetry and systematic absences of these photographs suggested the noncentrosymmetric orthorhombic space group  $P2_12_12_1$ .

The unit-cell dimensions were obtained from powder photographs recorded in a Guinier-Hägg XDC-700 focussing camera using  $\text{CuK}\alpha_1$  radiation ( $\lambda = 1.54051 \text{ \AA}$ ), with silicon ( $a = 5.43054 \text{ \AA}$  at  $25^\circ\text{C}$ ) as an internal standard. A least-squares treatment of the observed  $\theta$ -values from 63 reflexions provided the parameters given in Table I.\* The  $b$ -axis was found to be about twice the value reported by Shuvalov et al. (1).

### Structure Determination

The structure was solved by direct methods according to Karle and Karle (4). The normalized structure factor  $E_h$  was defined by  $E_h^2 = F_h^2/\epsilon \sum f_j^2$ , where  $\epsilon = 2$  when  $h$  is  $h00$ ,  $0k0$ , or  $00l$ , and  $\epsilon = 1$  otherwise;  $f_j$  is the atomic scattering factor for the  $j$ th atom in the unit cell containing  $N$  atoms. The set of reflexions used for the choice of origin, the enantiomorph and those given symbolic phases are shown in Table II.

The procedure for fixing the origin in a non-

\* Numbers in parentheses here and throughout this paper are the estimated standard deviations in the least-significant digits.

TABLE I  
CRYSTAL DATA<sup>a</sup>

Crystal system	Orthorhombic
Space group	$P2_12_12_1$
$a$	$5.9192(2) \text{ \AA}$
$b$	$17.9506(5) \text{ \AA}$
$c$	$6.2519(3) \text{ \AA}$
$V$	$664.287 \text{ \AA}^3$
$Z$	4
$D_x$	$3.42 \text{ g/cm}^3$
$\mu_{\text{calc}} (\text{MoK}\alpha)$	$189.7 \text{ cm}^{-1}$

<sup>a</sup>  $25^\circ\text{C}$ .

TABLE II  
STARTING SET PARAMETERS FOR PHASE DETERMINATION

$h$	$ E_h $	$\phi_h$	
1, 23, 0	3.03	$\pi/2$	} origin
3, 0, 4	1.76	0	
0, 11, 3	2.01	$\pi/2$	
0, 5, 8	1.70	$\pi/2$	} enantiomorph
1, 3, 1	2.82	$a$	} symbolic phases
4, 0, 2	2.53	$b$	
4, 20, 1	2.15	$c$	

centrosymmetric space group is discussed in detail by Hauptman and Karle (5). By using the formula  $\phi_h \approx \langle \phi_k + \phi_{h-k} \rangle_{k_r}$ , where  $\phi_h$  is the phase related to  $|E_h|$  and  $k_r$  means that only the phases corresponding to the largest  $|E_k|$  are considered, the phases of the symbols were found to be  $a \approx -\pi/2$ ,  $b = \pi$ , and  $c \approx 0$ . From these considerations 41 reflexions could be assigned phases and used as input to the tangent refinement,

$$\tan \phi_h \approx \frac{\sum |E_k \cdot E_{h-k}| \sin(\phi_k + \phi_{h-k})}{\sum |E_k \cdot E_{h-k}| \cos(\phi_k + \phi_{h-k})}$$

After a few cycles the phases of 121 reflexions with  $|E| > 1.49$  were accepted and an  $E$ -map could be calculated. This yielded the coordinates for the rubidium and selenium atoms and, after a least-squares refinement based on these three heavy atoms, a difference Fourier calculation gave the positions of the oxygen atoms.

### Refinement

The structure was refined by minimizing the function  $\sum w(|F_o| - |F_c|)^2$  using the full matrix least-squares program LINUS. Each reflexion was assigned a weight inversely proportional to the estimated variance of the observation  $w$ :  $w = 1/\sigma^2(F)$  where  $\sigma^2(F) = \sigma_{\text{count}}^2 + (F^2)/4F^2 + k \cdot F^2$ . The first term is based on counting statistics and the second is an empirical correction term to take account of other errors in the data. The constant  $k$  was chosen to be  $0.0004$ . The refinement was based on the 934 reflexions with intensities  $> 3\sigma_{\text{count}}$ . The parameters refined comprised 27 positional and 54 anisotropic thermal parameters, one overall scale factor, and one isotropic extinction parameter.

The positions of the hydrogen atoms could not be found in the difference Fourier maps. However, the Se-O distances clearly indicate to which oxygen atoms the three hydrogen atoms are

covalently bonded (cf. the "Description and Discussion ..." section below). The coordinates of the hydrogen positions were therefore calculated assuming linear O···O hydrogen bonds, with the atoms 1.05 Å from O(1), O(3), and O(6), respectively. They were included in the last cycles of calculation with fixed parameters. The isotropic thermal parameter for each was assumed to be 6.0 Å<sup>2</sup>.

The atomic scattering factors used for Se and O were those calculated by Hanson *et al.* (6), for Rb<sup>+</sup> those calculated by Cromer, Larson, and Waber (7), and for H as calculated by Stewart, Davidson, and Simpson (8). The scattering factors of Se and Rb<sup>+</sup> were corrected for anomalous dispersion according to Cromer (9). The final agreement factors were:

$$R = (\sum ||F_o| - |F_c|| / \sum |F_o|) = 0.030, \quad \text{and}$$

$$R_w = (\sum w(|F_o| - |F_c|)^2) / \sum wF_o^2)^{1/2} = 0.034.$$

The standard deviation of an observation of unit weight  $\sigma_1 = [\sum w(|F_o| - |F_c|)^2 / (m - n)]^{1/2}$  was 1.07, where  $m$  is the total number of observations and  $n$  the number of parameters varied. The final parameters are given in Tables III and IV, the observed and calculated structure factors in Table V. The observed  $F_o$  values are corrected for extinction using the expression

$$|F_o|_{\text{corr}} = |F_o| \left( 1 + \frac{\bar{T} \cdot p \cdot |F_c|^2 \cdot g \cdot \lambda^3}{12.593 \cdot V^2 \cdot \sin 2\theta} \right)^{1/4},$$

where  $\lambda$  is in Å,  $V$  is the unit cell volume in Å<sup>3</sup>, and  $p = 2 \cdot (1 + \cos^4 2\theta) / (1 + \cos^2 2\theta)$  (which is strictly valid only for unpolarized X-rays) (10). The mean path length through the crystal  $\bar{T}$  for a particular

TABLE III  
ATOMIC COORDINATES ( $\times 10^5$ )

	<i>x</i>	<i>y</i>	<i>z</i>
Se(1)	21630(14)	39527(4)	29576(12)
Se(2)	73113(14)	24215(4)	16878(12)
Rb	26034(15)	08289(4)	23398(13)
O(1)	18958(119)	42057(28)	02738(86)
O(2)	50635(98)	22483(34)	31855(138)
O(3)	49313(87)	42686(28)	34067(103)
O(4)	77185(96)	33470(24)	18904(88)
O(5)	08180(90)	46152(28)	42379(106)
O(6)	95210(96)	20750(31)	32267(145)

reflexion is appropriately corrected for absorption. The refined value of the isotropic extinction parameter is  $g = 4043(325)$ . This corresponds to a mosaic spread parameter of approximately 14 sec or a domain size of 0.3 nm, depending upon whether a Zachariasen (11) Type I or Type II description is assumed. The extinction correction factors as given above were in the range 1.00–1.25.

#### Absolute Configuration

Since the effect of anomalous dispersion was included in the structure factor calculation, a determination of the absolute configuration was possible (12). Refinement of a structure having an absolute configuration opposite to that in Table III converged with agreement factors  $R(-) = 0.036$  and  $R_w(-) = 0.041$ . Use of Hamilton's  $R$ -test (13) indicates a high significance for the absolute configuration given in Table III.

TABLE IV  
ANISOTROPIC THERMAL PARAMETERS ( $\times 10^5$ )<sup>a</sup>

	$\beta_{11}$	$\beta_{22}$	$\beta_{33}$	$\beta_{12}$	$\beta_{13}$	$\beta_{23}$	$R_1$	$R_2$	$R_3$
Se(1)	1418(21)	107(2)	1290(18)	-24(5)	-54(20)	-16(5)	131(1)	157(1)	162(1)
Se(2)	1546(23)	145(2)	1179(18)	-15(6)	199(21)	-38(5)	143(1)	155(1)	174(1)
Rb	1538(21)	169(2)	1541(19)	39(6)	70(22)	35(5)	159(1)	166(1)	180(1)
O(1)	3253(251)	216(18)	1135(132)	-242(57)	-408(166)	110(41)	140(9)	173(9)	257(9)
O(2)	1871(177)	281(22)	2787(249)	-243(51)	574(207)	-52(67)	152(10)	218(11)	252(11)
O(3)	1365(154)	153(16)	2242(183)	-1(40)	-70(165)	-51(46)	155(9)	157(8)	212(9)
O(4)	2093(169)	119(14)	2408(158)	14(44)	677(204)	42(40)	139(8)	170(11)	237(10)
O(5)	1483(162)	172(17)	2382(185)	60(45)	102(160)	-332(50)	123(11)	166(9)	242(9)
O(6)	1981(185)	216(20)	3111(263)	225(50)	-1023(207)	-113(66)	143(11)	192(10)	274(11)

<sup>a</sup> The form of the temperature factor is  $\exp[-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)]$ . The root mean square components  $R_i$  ( $\times 10^3$  Å) of thermal displacement of the atoms along the ellipsoid axes are also listed.

TABLE V

OBSERVED AND CALCULATED STRUCTURE FACTORS<sup>a, b</sup>

h	k	l	F <sub>o</sub>	F <sub>c</sub>
0	0	0	2777	2777
2	4	6	601	605
4	4	6	311	325
6	4	6	209	221
8	4	6	866	866
10	4	6	310	329
12	4	6	778	778
14	4	6	1173	1105
16	4	6	507	510
18	4	6	243	244
20	4	6	274	287
22	4	6	309	292
24	4	6	302	302
26	4	6	168	168
28	4	6	281	265
30	4	6	576	639
32	4	6	369	377
34	4	6	196	214
36	4	6	575	589
38	4	6	167	162
40	4	6	1304	1317
42	4	6	291	297
44	4	6	454	448
46	4	6	160	136
48	4	6	105	87
50	4	6	304	300
52	4	6	683	683
54	4	6	655	640
56	4	6	459	456
58	4	6	536	538
60	4	6	206	206
62	4	6	306	315
64	4	6	659	659
66	4	6	1013	1013
68	4	6	59	26
70	4	6	484	362
72	4	6	321	321
74	4	6	1230	1220
76	4	6	1835	1784
78	4	6	530	535
80	4	6	109	106
82	4	6	229	223
84	4	6	111	111
86	4	6	405	347
88	4	6	504	512
90	4	6	216	222
92	4	6	340	352
94	4	6	110	110
96	4	6	398	376
98	4	6	768	736
100	4	6	367	381
102	4	6	28	28
104	4	6	404	407
106	4	6	113	101
108	4	6	512	512
110	4	6	183	170
112	4	6	255	259
114	4	6	532	531
116	4	6	856	872
118	4	6	700	712
120	4	6	217	225
122	4	6	393	379
124	4	6	182	182
126	4	6	441	436
128	4	6	238	234
130	4	6	988	988
132	4	6	699	715
134	4	6	341	341
136	4	6	368	379
138	4	6	349	352
140	4	6	719	719
142	4	6	0	0
144	4	6	1095	1095
146	4	6	328	328
148	4	6	739	742
150	4	6	302	293
152	4	6	235	208
154	4	6	1347	1347
156	4	6	1108	1108
158	4	6	892	892
160	4	6	494	462
162	4	6	285	286
164	4	6	107	91
166	4	6	253	234
168	4	6	264	264
170	4	6	98	100
172	4	6	290	292
174	4	6	193	196
176	4	6	58	58
178	4	6	80	107
180	4	6	192	211
182	4	6	676	687
184	4	6	643	659
186	4	6	101	101
188	4	6	387	387
190	4	6	165	163
192	4	6	352	330
194	4	6	165	164
196	4	6	393	390
198	4	6	165	165
200	4	6	359	359
202	4	6	193	193
204	4	6	583	583
206	4	6	384	384
208	4	6	981	981
210	4	6	384	382
212	4	6	876	876
214	4	6	416	422
216	4	6	389	395
218	4	6	389	389
220	4	6	981	981
222	4	6	822	813
224	4	6	832	832
226	4	6	138	141
228	4	6	267	259
230	4	6	363	370
232	4	6	440	443
234	4	6	254	256
236	4	6	440	443
238	4	6	254	256
240	4	6	377	377
242	4	6	202	202
244	4	6	483	478
246	4	6	0	0
248	4	6	0	0
250	4	6	0	0
252	4	6	0	0
254	4	6	0	0
256	4	6	0	0
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620	4	6	0	0
622	4	6	0	0
624	4	6	0	0
626	4	6	0	0
628	4	6	0	0
630				



### Computer Programs

The structure was solved with an IBM 1800 version of the program PHASE written by Koenig (14). The final calculations were done on the CDC 3600 computer in Uppsala using programs briefly described by Jönsson and Liminga (15).

### Description and Discussion of the Structure

#### General

The structure is illustrated in Figs. 1–4. Bond distances and angles are listed in Table VI, with some distances also shown in the figures. The

atomic arrangement of the heavy atoms in the asymmetric unit is illustrated in Fig. 1. As was pointed out earlier, it was not possible to find the positions of the hydrogen atoms from the difference Fourier maps. However, as in the case of  $\text{LiH}_3(\text{SeO}_3)_2$  (2), the Se–O distances very clearly indicate that the hydrogen atoms are bonded to O(1), O(3), and O(6). From these considerations it is possible to formally describe the asymmetric unit as consisting of an  $\text{Rb}^+$  ion, an  $\text{HSeO}_3^-$  ion, and an  $\text{H}_2\text{SeO}_3$  molecule. The  $\text{HSe}(2)\text{O}_3^-$  ions are connected by hydrogen bonds  $\text{O}(2) \cdots \text{HO}(6)$ , forming S-shaped chains running in the [100]-direction (Fig. 2). The molecules  $\text{H}_2\text{Se}(1)\text{O}_3$  are

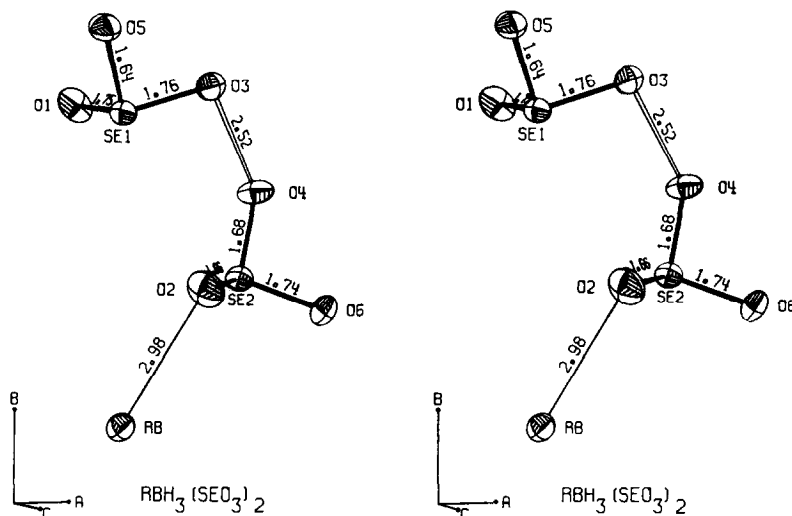


FIG. 1. A stereoscopic illustration of the heavy atoms of the asymmetric unit in  $\text{RbH}_3(\text{SeO}_3)_2$ . The ellipsoids are scaled to include 50% probability. Covalent bonds are filled,  $\text{O} \cdots \text{O}$  hydrogen bonds are open, and ionic bonds are drawn as single lines.

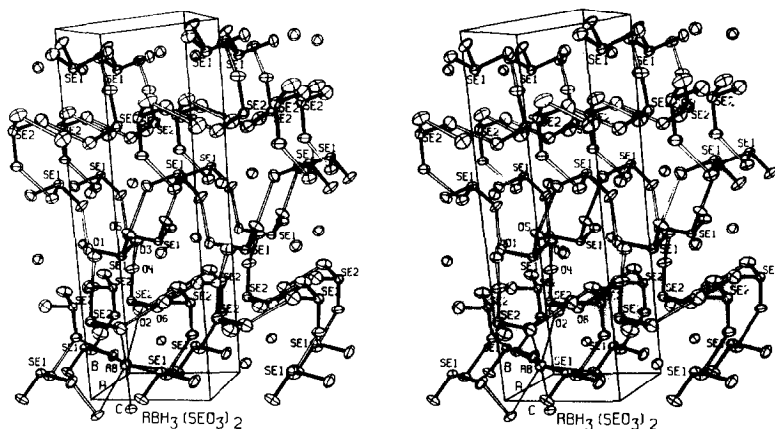


FIG. 2. A stereoscopic view of the structure showing the different types of chains.

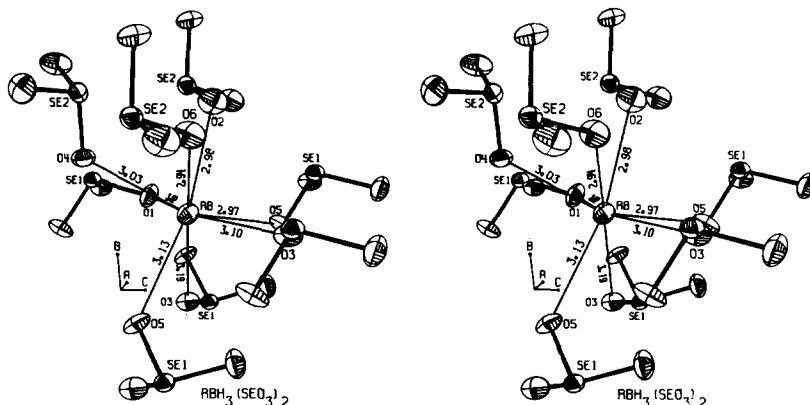


FIG. 3. A stereoscopic pair of drawings showing the rubidium ion and its surroundings.

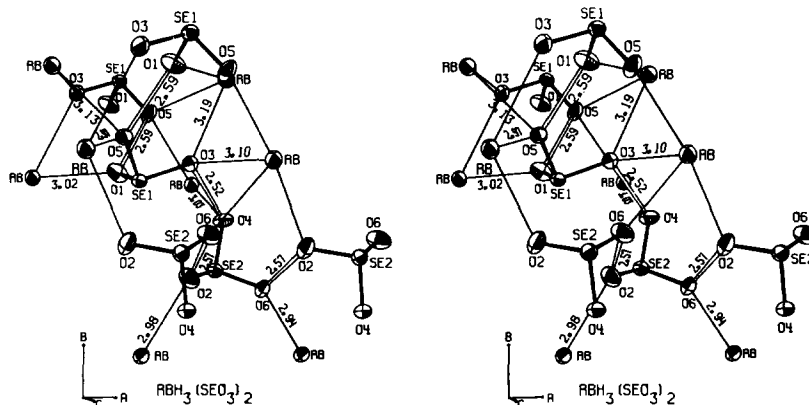


FIG. 4. A stereoscopic view of the hydrogen bond scheme.

also hydrogen bonded together via  $O(1)H \cdots O(5)$ , forming similar chains in the  $[001]$ -direction.

Puckered sheets of  $-Se(1)-Se(1)-$  and  $-Se(2)-Se(2)-$  chains with chain direction at right angles thus alternate in the  $[010]$ -direction. The different types of chains are interlinked through the third hydrogen bond  $O(3)H \cdots O(4)$  to form a three-dimensional network. The network is further stabilized by ionic bonds  $Rb^+-O$ .

#### The Rubidium Ion

The rubidium ion is surrounded by eight oxygen atoms forming a distorted cube (Fig. 3 and Table VI). The  $Rb^+-O$  distances are between 2.939(6) and 3.192(5) Å. As can be seen in Figs. 2 and 3,  $Rb^+ \cdots O$  bonds are formed with oxygen atoms situated both in the same and in adjacent chains, thus stabilizing the structure both within and between different chains.

#### The Selenium and Oxygen Atoms

There are two independent selenium atoms present in the structure. They are covalently

bonded to three oxygen atoms, each forming a pyramid with Se-O distances 1.640(6)–1.757(5) Å (Fig. 1 and Table VI). The angles O-Se-O are between  $98.09(26)^\circ$ , and  $104.93(29)^\circ$  (Table VI). Distances and angles are in good agreement with those in  $LiH_3(SeO_3)_2$  (2) and  $KH_3(SeO_3)_2$  (16). As in these cases, it is to be expected that the Se-O distances are dependent upon whether or not a hydrogen atom is covalently bonded to the oxygen atom. From consideration of the Se-O distances, it is clear that the hydrogen atoms are covalently bonded to O(1), O(3) and O(6), having Se-O distances 1.746(5), 1.757(5), and 1.739(7) Å, respectively. The oxygen atoms O(2), O(4), and O(5), with Se-O distances 1.656(7), 1.683(4), and 1.640(6) Å, respectively, thus serve as acceptor atoms in the hydrogen bonds.

#### The Hydrogen Bond System

The details of the hydrogen bond system are shown in Fig. 4. It consists of three different hydrogen bonds. The first  $O(1)H \cdots O(5)$ , 2.594(8) Å, connects the  $H_2Se(1)O_3$  groups

TABLE VI  
INTERATOMIC DISTANCES AND ANGLES IN  $\text{RbH}_3(\text{SeO}_3)_2^a$

A. Distances (Å)		B. Angles in degrees	
Se(1)–O(1)	1.746(5)	O(1)–Se(1)–O(3)	98.87(31)
–O(3)	1.757(5)	O(1)–Se(1)–O(5)	103.68(29)
–O(5)	1.640(6)	O(3)–Se(1)–O(5)	98.09(26)
Se(2)–O(2)	1.656(7)	O(2)–Se(2)–O(4)	104.93(29)
–O(4)	1.683(4)	O(2)–Se(2)–O(6)	102.96(36)
–O(6)	1.739(7)	O(4)–Se(2)–O(6)	101.76(27)
O(1)H···O(5) <sub>5642</sub>	2.594(8)	O(1) <sub>5553</sub> –Rb–O(2)	72.71(17)
O(6)H···O(2) <sub>5563</sub>	2.571(12)	O(1) <sub>5553</sub> –Rb–O(4) <sub>4553</sub>	61.27(15)
O(3)H···O(4)	2.521(7)	O(2)–Rb–O(6) <sub>4551</sub>	67.63(16)
Rb–O(1) <sub>5553</sub>	3.021(7)	O(4) <sub>4553</sub> –Rb–O(6) <sub>4551</sub>	78.88(19)
–O(2)	2.982(6)	O(3) <sub>4563</sub> –Rb–O(5) <sub>5454</sub>	84.32(15)
–O(3) <sub>4563</sub>	3.099(6)	O(3) <sub>4563</sub> –Rb–O(5) <sub>5563</sub>	72.17(15)
–O(3) <sub>6454</sub>	3.192(5)	O(3) <sub>6454</sub> –Rb–O(5) <sub>5454</sub>	68.87(13)
–O(4) <sub>4553</sub>	3.031(5)	O(3) <sub>6454</sub> –Rb–O(5) <sub>5563</sub>	65.00(15)
–O(5) <sub>5563</sub>	2.972(6)	O(1) <sub>5553</sub> –Rb–O(3) <sub>6454</sub>	61.22(14)
–O(5) <sub>5454</sub>	3.134(5)	O(2)–Rb–O(5) <sub>5563</sub>	77.81(17)
–O(6) <sub>4551</sub>	2.939(6)	O(4) <sub>4553</sub> –Rb–O(5) <sub>5454</sub>	94.54(15)
		O(6) <sub>4551</sub> –Rb–O(3) <sub>4563</sub>	64.17(19)

<sup>a</sup> Atoms outside the asymmetric unit are specified by a subscript which indicates how the atomic parameters can be derived from those in the asymmetric unit. The first three digits code the lattice translation (the asymmetric unit is in cell 555) and the fourth digit specifies the symmetry operation: 1 =  $x, y, z$ ; 2 =  $1/2 - x, -y, 1/2 + z$ ; 3 =  $1/2 + x, 1/2 - y, -z$ ; 4 =  $x, 1/2 + y, 1/2 - z$ .

to form S-shaped chains and the second bond O(2)···HO(6), 2.571(12) Å, connects the  $\text{HSe(2)O}_3^-$  groups in a similar way. The third bond, O(3)H···O(4) 2.521(7) Å, links the Se(1)– and Se(2)–chains. This scheme is completely different from that found in  $\text{LiH}_3(\text{SeO}_3)_2$  (2), but to discuss it in detail it is necessary to know the accurate hydrogen positions.

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