

Hydrogen Bond Studies

67.* The Crystal Structure of Rubidium Trihydrogen Selenite, $\text{RbH}_3(\text{SeO}_3)_2$

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The crystal structure of $\text{RbH}_3(\text{SeO}_3)_2$ has been determined from three-dimensional single crystal X-ray diffractometer data obtained at room temperature. Four formula units crystallize in an orthorhombic unit cell of dimensions: $a = 5.9192(2)$, $b = 17.9506(5)$, and $c = 6.2519(3)$ Å. The space group is $P2_1 2_1 2_1$. The structure consists of two types of chains at a right angle. One chain is built up of H_2SeO_3 molecules linked by 2.594(8)-Å hydrogen bonds and the other of HSeO_3^- ions linked by 2.571(12)-Å hydrogen bonds. These two types of chains are cross-linked by a third hydrogen bond of length 2.521(7) Å. The rubidium ion is surrounded by eight oxygen atoms forming a distorted cube. The Rb^+-O distances are in the range 2.94–3.19 Å.

Introduction

Crystals of the alkali trihydrogen selenites, $\text{MH}_3(\text{SeO}_3)_2$, have been found to show a variety of dielectric properties and crystal structures [for a review see Ref. (1)]. To be able to explain these different properties it is necessary to have detailed structural information.

The present investigation forms part of a series of X-ray and neutron diffraction studies of hydrogen-bonded ferroelectric and related compounds in progress at this Institute. The previous paper was a neutron diffraction study of $\text{LiH}_3(\text{SeO}_3)_2$ (2), and the next will be an X-ray diffraction study of $\text{NH}_4\text{H}_3(\text{SeO}_3)_2$ (3).

Experimental

Crystals of $\text{RbH}_3(\text{SeO}_3)_2$ were prepared from an aqueous solution of Rb_2CO_3 and SeO_2 in the stoichiometric ratio. The solution was heated to 80°C. On cooling to room temperature small almost spherical single crystals were obtained. Their melting point was found to be 57°C, which agrees very well with the reported value of 57.1°C (1).

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The intensity data were collected at room temperature on a Stoe-Philips computer controlled four-circle diffractometer with a graphite monochromator using $\text{MoK}\alpha$ radiation and $\omega-2\theta$ scan. Different combinations of scan time and filter selection were used with the intention of measuring all reflections to about the same precision.

Three standard reflexions with integrated intensities of 3500, 4000, and 47 000 counts, respectively, were measured at regular intervals to provide a check on crystal and electronic stability. The two weaker reflexions showed no significant changes, whereas the intensity of the strongest one first underwent a continuous increase of 10% and then stabilized. All reflexions obtained up to this point were later remeasured. It was then found that the increase had affected only the high-intensity reflexions. This increase was probably due to an alteration in the nature of the extinction effects present. It could be that the mosaic structure had been changed because of radiation damage. As the data used in the refinements were those collected after the test reflexion had stabilized, no scaling was necessary.

A total of 1172 independent reflexions having $\sin \theta/\lambda$ less than 0.74 \AA^{-1} were recorded. Out of these, 934 had intensities greater than $3\sigma_{\text{count}}$. The intensities were corrected for Lorentz, polarization, and absorption effects (spherical

crystal). The calculated linear absorption coefficient was 189.7 cm^{-1} . The crystal was approximately spherical with a radius of 0.0376 mm . The transmission factors were in the range $0.356\text{--}0.373$.

Space Group and Unit Cell

The space group was determined from Weissenberg photographs recorded with $\text{CuK}\alpha$ radiation. The diffraction symmetry and systematic absences of these photographs suggested the noncentrosymmetric orthorhombic space group $P2_12_12_1$.

The unit-cell dimensions were obtained from powder photographs recorded in a Guinier-Hägg XDC-700 focussing camera using $\text{CuK}\alpha_1$ radiation ($\lambda = 1.54051 \text{ \AA}$), with silicon ($a = 5.43054 \text{ \AA}$ at 25°C) as an internal standard. A least-squares treatment of the observed θ -values from 63 reflexions provided the parameters given in Table I.* The b -axis was found to be about twice the value reported by Shuvalov et al. (1).

Structure Determination

The structure was solved by direct methods according to Karle and Karle (4). The normalized structure factor E_h was defined by $E_h^2 = F_h^2/\epsilon \sum f_j^2$, where $\epsilon = 2$ when h is $h00$, $0k0$, or $00l$, and $\epsilon = 1$ otherwise; f_j is the atomic scattering factor for the j th atom in the unit cell containing N atoms. The set of reflexions used for the choice of origin, the enantiomorph and those given symbolic phases are shown in Table II.

The procedure for fixing the origin in a non-

* Numbers in parentheses here and throughout this paper are the estimated standard deviations in the least-significant digits.

TABLE I
CRYSTAL DATA^a

Crystal system	Orthorhombic
Space group	$P2_12_12_1$
a	$5.9192(2) \text{ \AA}$
b	$17.9506(5) \text{ \AA}$
c	$6.2519(3) \text{ \AA}$
V	664.287 \AA^3
Z	4
D_x	3.42 g/cm^3
$\mu_{\text{calc}} (\text{MoK}\alpha)$	189.7 cm^{-1}

^a 25°C .

TABLE II
STARTING SET PARAMETERS FOR PHASE DETERMINATION

h	$ E_h $	ϕ_h	
1, 23, 0	3.03	$\pi/2$	} origin
3, 0, 4	1.76	0	
0, 11, 3	2.01	$\pi/2$	
0, 5, 8	1.70	$\pi/2$	} enantiomorph
1, 3, 1	2.82	a	} symbolic phases
4, 0, 2	2.53	b	
4, 20, 1	2.15	c	

centrosymmetric space group is discussed in detail by Hauptman and Karle (5). By using the formula $\phi_h \approx \langle \phi_k + \phi_{h-k} \rangle_{k_r}$, where ϕ_h is the phase related to $|E_h|$ and k_r means that only the phases corresponding to the largest $|E_k|$ are considered, the phases of the symbols were found to be $a \approx -\pi/2$, $b = \pi$, and $c \approx 0$. From these considerations 41 reflexions could be assigned phases and used as input to the tangent refinement,

$$\tan \phi_h \approx \frac{\sum |E_k \cdot E_{h-k}| \sin(\phi_k + \phi_{h-k})}{\sum |E_k \cdot E_{h-k}| \cos(\phi_k + \phi_{h-k})}$$

After a few cycles the phases of 121 reflexions with $|E| > 1.49$ were accepted and an E -map could be calculated. This yielded the coordinates for the rubidium and selenium atoms and, after a least-squares refinement based on these three heavy atoms, a difference Fourier calculation gave the positions of the oxygen atoms.

Refinement

The structure was refined by minimizing the function $\sum w(|F_o| - |F_c|)^2$ using the full matrix least-squares program LINUS. Each reflexion was assigned a weight inversely proportional to the estimated variance of the observation w : $w = 1/\sigma^2(F)$ where $\sigma^2(F) = \sigma_{\text{count}}^2 + (F^2)/4F^2 + k \cdot F^2$. The first term is based on counting statistics and the second is an empirical correction term to take account of other errors in the data. The constant k was chosen to be 0.0004 . The refinement was based on the 934 reflexions with intensities $> 3\sigma_{\text{count}}$. The parameters refined comprised 27 positional and 54 anisotropic thermal parameters, one overall scale factor, and one isotropic extinction parameter.

The positions of the hydrogen atoms could not be found in the difference Fourier maps. However, the Se-O distances clearly indicate to which oxygen atoms the three hydrogen atoms are

covalently bonded (cf. the "Description and Discussion ..." section below). The coordinates of the hydrogen positions were therefore calculated assuming linear O···O hydrogen bonds, with the atoms 1.05 Å from O(1), O(3), and O(6), respectively. They were included in the last cycles of calculation with fixed parameters. The isotropic thermal parameter for each was assumed to be 6.0 Å².

The atomic scattering factors used for Se and O were those calculated by Hanson *et al.* (6), for Rb⁺ those calculated by Cromer, Larson, and Waber (7), and for H as calculated by Stewart, Davidson, and Simpson (8). The scattering factors of Se and Rb⁺ were corrected for anomalous dispersion according to Cromer (9). The final agreement factors were:

$$R = (\sum ||F_o| - |F_c|| / \sum |F_o|) = 0.030, \text{ and}$$

$$R_w = (\sum w(|F_o| - |F_c|)^2) / \sum wF_o^2)^{1/2} = 0.034.$$

The standard deviation of an observation of unit weight $\sigma_1 = [\sum w(|F_o| - |F_c|)^2 / (m - n)]^{1/2}$ was 1.07, where m is the total number of observations and n the number of parameters varied. The final parameters are given in Tables III and IV, the observed and calculated structure factors in Table V. The observed F_o values are corrected for extinction using the expression

$$|F_o|_{\text{corr}} = |F_o| \left(1 + \frac{\bar{T} \cdot p \cdot |F_c|^2 \cdot g \cdot \lambda^3}{12.593 \cdot V^2 \cdot \sin 2\theta} \right)^{1/4},$$

where λ is in Å, V is the unit cell volume in Å³, and $p = 2 \cdot (1 + \cos^4 2\theta) / (1 + \cos^2 2\theta)$ (which is strictly valid only for unpolarized X-rays) (10). The mean path length through the crystal \bar{T} for a particular

TABLE III
ATOMIC COORDINATES ($\times 10^5$)

	<i>x</i>	<i>y</i>	<i>z</i>
Se(1)	21630(14)	39527(4)	29576(12)
Se(2)	73113(14)	24215(4)	16878(12)
Rb	26034(15)	08289(4)	23398(13)
O(1)	18958(119)	42057(28)	02738(86)
O(2)	50635(98)	22483(34)	31855(138)
O(3)	49313(87)	42686(28)	34067(103)
O(4)	77185(96)	33470(24)	18904(88)
O(5)	08180(90)	46152(28)	42379(106)
O(6)	95210(96)	20750(31)	32267(145)

reflexion is appropriately corrected for absorption. The refined value of the isotropic extinction parameter is $g = 4043(325)$. This corresponds to a mosaic spread parameter of approximately 14 sec or a domain size of 0.3 nm, depending upon whether a Zachariasen (11) Type I or Type II description is assumed. The extinction correction factors as given above were in the range 1.00–1.25.

Absolute Configuration

Since the effect of anomalous dispersion was included in the structure factor calculation, a determination of the absolute configuration was possible (12). Refinement of a structure having an absolute configuration opposite to that in Table III converged with agreement factors $R(-) = 0.036$ and $R_w(-) = 0.041$. Use of Hamilton's R -test (13) indicates a high significance for the absolute configuration given in Table III.

TABLE IV
ANISOTROPIC THERMAL PARAMETERS ($\times 10^5$)^a

	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}	R_1	R_2	R_3
Se(1)	1418(21)	107(2)	1290(18)	-24(5)	-54(20)	-16(5)	131(1)	157(1)	162(1)
Se(2)	1546(23)	145(2)	1179(18)	-15(6)	199(21)	-38(5)	143(1)	155(1)	174(1)
Rb	1538(21)	169(2)	1541(19)	39(6)	70(22)	35(5)	159(1)	166(1)	180(1)
O(1)	3253(251)	216(18)	1135(132)	-242(57)	-408(166)	110(41)	140(9)	173(9)	257(9)
O(2)	1871(177)	281(22)	2787(249)	-243(51)	574(207)	-52(67)	152(10)	218(11)	252(11)
O(3)	1365(154)	153(16)	2242(183)	-1(40)	-70(165)	-51(46)	155(9)	157(8)	212(9)
O(4)	2093(169)	119(14)	2408(158)	14(44)	677(204)	42(40)	139(8)	170(11)	237(10)
O(5)	1483(162)	172(17)	2382(185)	60(45)	102(160)	-332(50)	123(11)	166(9)	242(9)
O(6)	1981(185)	216(20)	3111(263)	225(50)	-1023(207)	-113(66)	143(11)	192(10)	274(11)

^a The form of the temperature factor is $\exp[-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)]$. The root mean square components R_i ($\times 10^3$ Å) of thermal displacement of the atoms along the ellipsoid axes are also listed.

TABLE V

OBSERVED AND CALCULATED STRUCTURE FACTORS^{a, b}

h	k	l	F _o	F _c	F _c /F _o
0	0	0	2777	2777	1.000
0	0	1	605	605	1.000
0	0	2	321	321	1.000
0	0	3	170	170	1.000
0	0	4	126	126	1.000
0	0	5	92	92	1.000
0	0	6	56	56	1.000
0	0	7	31	31	1.000
0	0	8	18	18	1.000
0	0	9	10	10	1.000
0	0	10	6	6	1.000
0	0	11	4	4	1.000
0	0	12	3	3	1.000
0	0	13	2	2	1.000
0	0	14	1	1	1.000
0	0	15	0	0	1.000
0	1	0	105	105	1.000
0	1	1	170	170	1.000
0	1	2	209	209	1.000
0	1	3	229	229	1.000
0	1	4	243	243	1.000
0	1	5	254	254	1.000
0	1	6	263	263	1.000
0	1	7	270	270	1.000
0	1	8	277	277	1.000
0	1	9	281	281	1.000
0	1	10	285	285	1.000
0	1	11	288	288	1.000
0	1	12	290	290	1.000
0	1	13	291	291	1.000
0	1	14	292	292	1.000
0	1	15	293	293	1.000
0	1	16	294	294	1.000
0	1	17	294	294	1.000
0	1	18	294	294	1.000
0	1	19	294	294	1.000
0	1	20	294	294	1.000
0	1	21	294	294	1.000
0	1	22	294	294	1.000
0	1	23	294	294	1.000
0	1	24	294	294	1.000
0	1	25	294	294	1.000
0	1	26	294	294	1.000
0	1	27	294	294	1.000
0	1	28	294	294	1.000
0	1	29	294	294	1.000
0	1	30	294	294	1.000
0	1	31	294	294	1.000
0	1	32	294	294	1.000
0	1	33	294	294	1.000
0	1	34	294	294	1.000
0	1	35	294	294	1.000
0	1	36	294	294	1.000
0	1	37	294	294	1.000
0	1	38	294	294	1.000
0	1	39	294	294	1.000
0	1	40	294	294	1.000
0	1	41	294	294	1.000
0	1	42	294	294	1.000
0	1	43	294	294	1.000
0	1	44	294	294	1.000
0	1	45	294	294	1.000
0	1	46	294	294	1.000
0	1	47	294	294	1.000
0	1	48	294	294	1.000
0	1	49	294	294	1.000
0	1	50	294	294	1.000
0	1	51	294	294	1.000
0	1	52	294	294	1.000
0	1	53	294	294	1.000
0	1	54	294	294	1.000
0	1	55	294	294	1.000
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0	1	68	294	294	1.000
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0	1	218	294	294	1.000
0	1	219	294	294	1.000
0	1	220	294	294	1.000
0	1	221	294	294	1.000
0	1	222	294	294	1.000
0	1	223	294	294	1.000
0	1				

Computer Programs

The structure was solved with an IBM 1800 version of the program PHASE written by Koenig (14). The final calculations were done on the CDC 3600 computer in Uppsala using programs briefly described by Jönsson and Liminga (15).

Description and Discussion of the Structure

General

The structure is illustrated in Figs. 1–4. Bond distances and angles are listed in Table VI, with some distances also shown in the figures. The

atomic arrangement of the heavy atoms in the asymmetric unit is illustrated in Fig. 1. As was pointed out earlier, it was not possible to find the positions of the hydrogen atoms from the difference Fourier maps. However, as in the case of $\text{LiH}_3(\text{SeO}_3)_2$ (2), the Se–O distances very clearly indicate that the hydrogen atoms are bonded to O(1), O(3), and O(6). From these considerations it is possible to formally describe the asymmetric unit as consisting of an Rb^+ ion, an HSeO_3^- ion, and an H_2SeO_3 molecule. The $\text{HSe}(2)\text{O}_3^-$ ions are connected by hydrogen bonds $\text{O}(2) \cdots \text{HO}(6)$, forming S-shaped chains running in the [100]-direction (Fig. 2). The molecules $\text{H}_2\text{Se}(1)\text{O}_3$ are

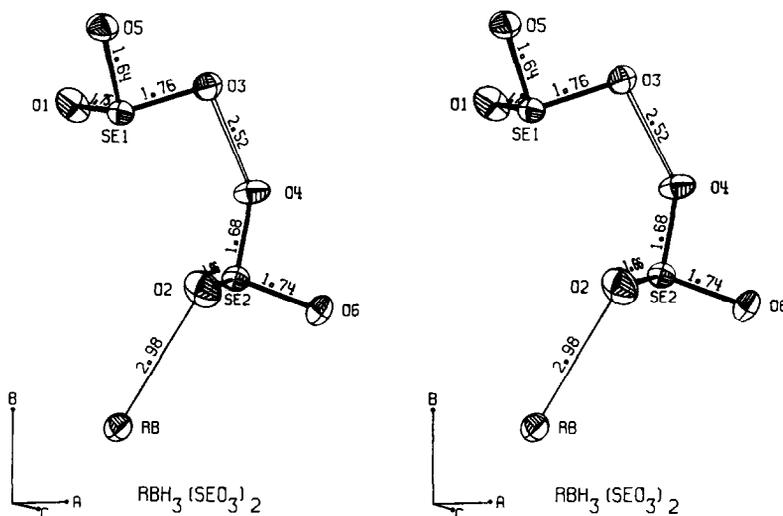


FIG. 1. A stereoscopic illustration of the heavy atoms of the asymmetric unit in $\text{RbH}_3(\text{SeO}_3)_2$. The ellipsoids are scaled to include 50% probability. Covalent bonds are filled, $\text{O} \cdots \text{O}$ hydrogen bonds are open, and ionic bonds are drawn as single lines.

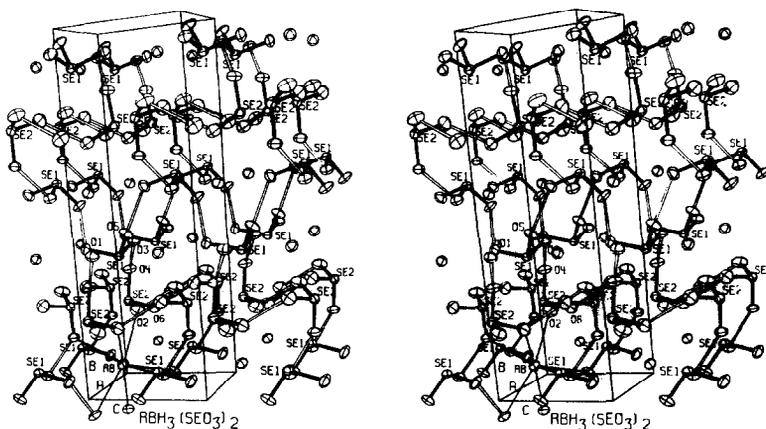


FIG. 2. A stereoscopic view of the structure showing the different types of chains.

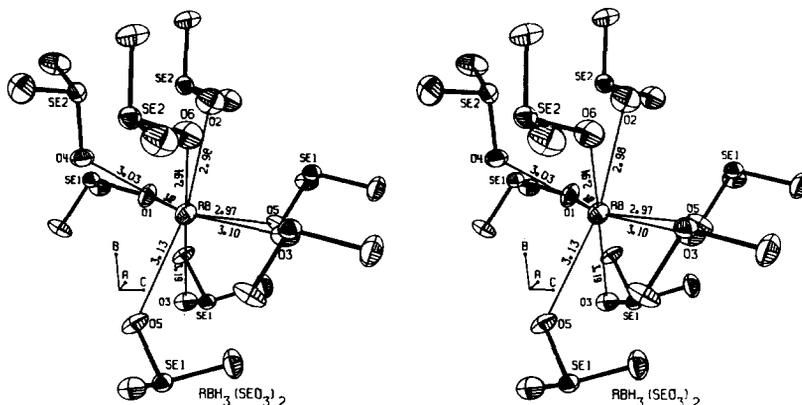


FIG. 3. A stereoscopic pair of drawings showing the rubidium ion and its surroundings.

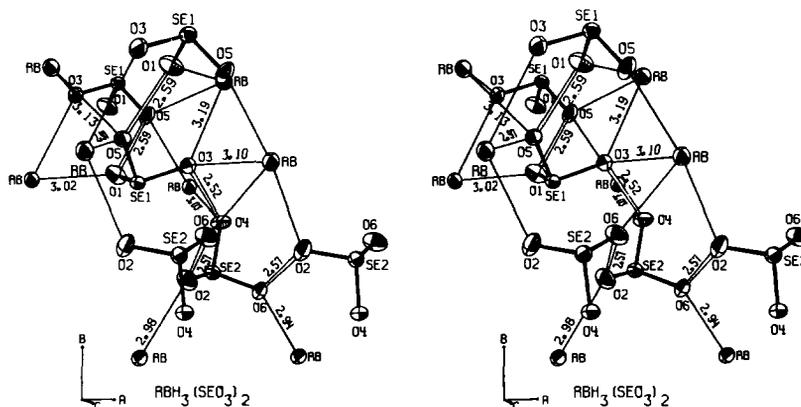


FIG. 4. A stereoscopic view of the hydrogen bond scheme.

also hydrogen bonded together via $O(1)H \cdots O(5)$, forming similar chains in the $[001]$ -direction.

Puckered sheets of $-Se(1)-Se(1)-$ and $-Se(2)-Se(2)-$ chains with chain direction at right angles thus alternate in the $[010]$ -direction. The different types of chains are interlinked through the third hydrogen bond $O(3)H \cdots O(4)$ to form a three-dimensional network. The network is further stabilized by ionic bonds Rb^+-O .

The Rubidium Ion

The rubidium ion is surrounded by eight oxygen atoms forming a distorted cube (Fig. 3 and Table VI). The Rb^+-O distances are between 2.939(6) and 3.192(5) Å. As can be seen in Figs. 2 and 3, $Rb^+ \cdots O$ bonds are formed with oxygen atoms situated both in the same and in adjacent chains, thus stabilizing the structure both within and between different chains.

The Selenium and Oxygen Atoms

There are two independent selenium atoms present in the structure. They are covalently

bonded to three oxygen atoms, each forming a pyramid with Se-O distances 1.640(6)–1.757(5) Å (Fig. 1 and Table VI). The angles O-Se-O are between $98.09(26)^\circ$, and $104.93(29)^\circ$ (Table VI). Distances and angles are in good agreement with those in $LiH_3(SeO_3)_2$ (2) and $KH_3(SeO_3)_2$ (16). As in these cases, it is to be expected that the Se-O distances are dependent upon whether or not a hydrogen atom is covalently bonded to the oxygen atom. From consideration of the Se-O distances, it is clear that the hydrogen atoms are covalently bonded to O(1), O(3) and O(6), having Se-O distances 1.746(5), 1.757(5), and 1.739(7) Å, respectively. The oxygen atoms O(2), O(4), and O(5), with Se-O distances 1.656(7), 1.683(4), and 1.640(6) Å, respectively, thus serve as acceptor atoms in the hydrogen bonds.

The Hydrogen Bond System

The details of the hydrogen bond system are shown in Fig. 4. It consists of three different hydrogen bonds. The first $O(1)H \cdots O(5)$, 2.594(8) Å, connects the $H_2Se(1)O_3$ groups

TABLE VI
INTERATOMIC DISTANCES AND ANGLES IN $\text{RbH}_3(\text{SeO}_3)_2^a$

A. Distances (Å)		B. Angles in degrees	
Se(1)–O(1)	1.746(5)	O(1)–Se(1)–O(3)	98.87(31)
–O(3)	1.757(5)	O(1)–Se(1)–O(5)	103.68(29)
–O(5)	1.640(6)	O(3)–Se(1)–O(5)	98.09(26)
Se(2)–O(2)	1.656(7)	O(2)–Se(2)–O(4)	104.93(29)
–O(4)	1.683(4)	O(2)–Se(2)–O(6)	102.96(36)
–O(6)	1.739(7)	O(4)–Se(2)–O(6)	101.76(27)
O(1)H···O(5) ₅₆₄₂	2.594(8)	O(1) ₅₅₅₃ –Rb–O(2)	72.71(17)
O(6)H···O(2) ₅₅₆₃	2.571(12)	O(1) ₅₅₅₃ –Rb–O(4) ₄₅₅₃	61.27(15)
O(3)H···O(4)	2.521(7)	O(2)–Rb–O(6) ₄₅₅₁	67.63(16)
Rb–O(1) ₅₅₅₃	3.021(7)	O(4) ₄₅₅₃ –Rb–O(6) ₄₅₅₁	78.88(19)
–O(2)	2.982(6)	O(3) ₄₅₆₃ –Rb–O(5) ₅₄₅₄	84.32(15)
–O(3) ₄₅₆₃	3.099(6)	O(3) ₄₅₆₃ –Rb–O(5) ₅₅₆₃	72.17(15)
–O(3) ₆₄₅₄	3.192(5)	O(3) ₆₄₅₄ –Rb–O(5) ₅₄₅₄	68.87(13)
–O(4) ₄₅₅₃	3.031(5)	O(3) ₆₄₅₄ –Rb–O(5) ₅₅₆₃	65.00(15)
–O(5) ₅₅₆₃	2.972(6)	O(1) ₅₅₅₃ –Rb–O(3) ₆₄₅₄	61.22(14)
–O(5) ₅₄₅₄	3.134(5)	O(2)–Rb–O(5) ₅₅₆₃	77.81(17)
–O(6) ₄₅₅₁	2.939(6)	O(4) ₄₅₅₃ –Rb–O(5) ₅₄₅₄	94.54(15)
		O(6) ₄₅₅₁ –Rb–O(3) ₄₅₆₃	64.17(19)

^a Atoms outside the asymmetric unit are specified by a subscript which indicates how the atomic parameters can be derived from those in the asymmetric unit. The first three digits code the lattice translation (the asymmetric unit is in cell 555) and the fourth digit specifies the symmetry operation: 1 = x, y, z ; 2 = $1/2 - x, -y, 1/2 + z$; 3 = $1/2 + x, 1/2 - y, -z$; 4 = $x, 1/2 + y, 1/2 - z$.

to form S-shaped chains and the second bond O(2)···HO(6), 2.571(12) Å, connects the HSe(2)O_3^- groups in a similar way. The third bond, O(3)H···O(4) 2.521(7) Å, links the Se(1)– and Se(2)–chains. This scheme is completely different from that found in $\text{LiH}_3(\text{SeO}_3)_2$ (2), but to discuss it in detail it is necessary to know the accurate hydrogen positions.

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References

1. L. A. SHUVALOV, N. R. IVANOV, N. V. GORDEEVA, AND L. F. KIRPICHNIKOVA, *Sov. Phys.-Crystallogr.* **14**, 554 (1970).
2. R. TELLGREN AND R. LIMINGA, *J. Solid State Chem.* **4**, 255 (1972).
3. R. TELLGREN, D. AHMAD, AND R. LIMINGA, *Chemica Scripta* in press.
4. J. KARLE AND I. L. KARLE, *Acta Crystallogr.* **21**, 849 (1966).
5. H. HAUPTMAN AND J. KARLE, *Acta Crystallogr.* **9**, 45 (1956).
6. H. P. HANSON, F. HERMAN, J. D. LEA, AND S. SKILLMAN, *Acta Crystallogr.* **17**, 1040 (1964).
7. D. T. CROMER, A. C. LARSON, AND J. T. WABER, *Acta Crystallogr.* **17**, 1044 (1964).
8. R. F. STEWART, E. R. DAVIDSON, AND W. T. SIMPSON, *J. Chem. Phys.* **42**, 3175 (1965).
9. D. T. CROMER, *Acta Crystallogr.* **18**, 17 (1965).
10. P. COPPENS AND W. C. HAMILTON, *Acta Crystallogr. Sect. A* **26**, 71 (1970).
11. W. H. ZACHARIASEN, *Acta Crystallogr.* **23**, 558 (1967).
12. J. A. IBERS AND W. C. HAMILTON, *Acta Crystallogr.* **17**, 781 (1964).
13. W. C. HAMILTON, *Acta Crystallogr.* **18**, 502 (1965).
14. D. F. KOENIG, Private communication.
15. P.-G. JÖNSSON AND R. LIMINGA, *Acta Chem. Scand.* **25**, 1729 (1971).
16. F. HANSEN, R. GRØNBÆK HAZELL, AND S. E. RASMUSSEN, *Acta Chem. Scand.* **23**, 2561 (1969).